

Magnetic fields and chemical peculiarities of the very young intermediate-mass binary system HD 72106

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ABSTRACT

The recently discovered magnetic Herbig Ae and Be stars may provide qualitatively new information about the formation and evolution of magnetic Ap and Bp stars. We have performed a detailed investigation of one particularly interesting binary system with a Herbig Ae secondary and a late B-type primary possessing a strong, globally ordered magnetic field. 20 high-resolution Stokes *V* spectra of the system were obtained with the ESPaDO*n*S instrument mounted on the Canada–France–Hawaii Telescope. In these observations we see clear evidence for a magnetic field in the primary, but no evidence for a magnetic field in the secondary. A detailed abundance analysis was performed for both stars, revealing strong chemical peculiarities in the primary and normal chemical abundances in the secondary. The primary is strongly overabundant in Si, Cr and other iron-peak elements, as well as Nd, and underabundant in He. The primary therefore appears to be a very young Bp star. In this context, line profile variations of the primary suggest non-uniform lateral distributions of surface abundances. Interpreting the 0.63995 ± 0.00009 d variation period of the Stokes *I* and *V* profiles as the rotational period of the star, we have modelled the magnetic field geometry and the surface abundance distributions of Si, Ti, Cr and Fe using magnetic Doppler imaging. We derive a dipolar geometry of the surface magnetic field, with a polar strength $B_d = 1230$ G and an obliquity $\beta = 57^\circ$. The distributions Ti, Cr and Fe are all qualitatively similar, with an elongated patch of enhanced abundance situated near the positive magnetic pole. The Si distribution is somewhat different, and its relationship to the magnetic field geometry less clear.

Key words: stars: abundances – stars: chemically peculiar – stars: evolution – stars: individual: HD 72106 – stars: magnetic fields.

1 INTRODUCTION

Strong, globally organized magnetic fields have recently been reported in a few Herbig Ae and Be (HAeBe) stars (Donati et al. 1997; Hubrig, Schöller & Yudin 2004; Wade et al. 2005; Catala et al. 2007; Wade et al. 2007; Alecian et al. 2008a,b). HAeBe stars

are pre-main-sequence stars of intermediate-mass which evolve to become main-sequence A and B stars. HAeBe stars have A or B spectral classes, display emission lines and infrared excesses, and are usually found with nebulosity nearby (Vieira et al. 2003). The detection of magnetic fields in HAeBe stars is of particular interest as it hints at a connection to the main-sequence magnetic, chemically peculiar Ap and Bp stars.

Ap and Bp stars display strong magnetic fields, with typical strengths of a few hundred gauss up to a few tens of kilogauss,

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and globally ordered geometries of approximately dipolar topology. These stars also display strong and distinctive chemical peculiarities, particularly overabundances of Si and iron peak elements, occasionally in excess of 2 dex, and even greater overabundances of some rare earth elements.

The source of the magnetic fields observed in Ap/Bp stars is not well understood. The two major competing field origin theories propose, on the one hand, that the field is a relic of an earlier stage of stellar evolution, now frozen into the plasma of the star or, on the other hand, that the field is generated contemporaneously by a dynamo. Also unknown are the details of the formation of the observed chemical peculiarities. Chemical peculiarities are believed to be the result of atomic diffusion leading to a chemically stratified stellar envelope (Michaud 1970; Michaud, Charland & Megessier 1981). However, the details of this process, particularly for individual elements, and the impact of magnetic fields are not fully understood (e.g. Alecian & Stift 2007). In this context it is of great interest to identify and characterize the evolutionary progenitors of Ap/Bp stars, as these objects can provide critical information on both the origin of magnetic fields and the formation of chemical peculiarities.

The recently discovered magnetic HAeBe stars have been proposed to be pre-main-sequence progenitors of Ap/Bp stars (Wade et al. 2005). Thus a detailed investigation of these stars can shed light on these questions, as well as potentially provide some information about magnetic braking and magnetospheric accretion in intermediate-mass stars.

2 HD 72106

HD 72106 is a double-star system in the constellation Vela, with a 0.805 arcsec separation between components (ESA 1997). The brighter star, HD 72106A ('the primary'), was identified as a magnetic star by Wade et al. (2005), while the fainter star, HD 72106B ('the secondary'), was identified as a HAeBe star by Vieira et al. (2003). The HD 72106 system was observed by *Hipparcos* (ESA 1997), and included in the recent re-reduction of *Hipparcos* data by van Leeuwen (2007a,b) who found a parallax of 3.60 ± 14 mas, placing the system at a heliocentric distance of 278_{-67}^{+129} pc. The system was identified as having an infrared excess, based on the *Infrared Astronomical Satellite* (IRAS) data, by Oudmaijer et al. (1992). Torres et al. (1995) observed HD 72106 and noted the presence of emission in the H α Balmer line of the combined spectrum of the system. Schütz, Meeus & Sterzik (2005) obtained infrared spectroscopy of HD 72106 system and examined circumstellar abundances of a number of dust grain species.

Vieira et al. (2003) associated the system with the Gum nebula star-forming region. They observed weak H α emission and a small contribution from dust in the spectral energy distribution of HD 72106B. This, they hypothesized, was due to HD 72106B being an evolved HAeBe star that has cleared most of its circumstellar envelope.

The discovery of a magnetic field in HD 72106A by Wade et al. (2005) was based on spectropolarimetry from FORS1 (FOcal Reducer/low dispersion Spectrograph) at the Very Large Telescope and from ESPaDOnS (Echelle SpectroPolarimetric Device for the Observation of Stars) at the Canada–France–Hawaii Telescope (CFHT). A longitudinal field of 195 ± 40 G (i.e. 4.9σ) was deduced from the FORS1 spectrum. No significant magnetic field was detected in the secondary (65 ± 55 G was observed). Although the average longitudinal field in the ESPaDOnS spectrum of the primary was consistent with zero, a clear circular polariza-

tion signature was detected with a very high degree of significance, unambiguously indicating the presence of a surface magnetic field.

A more conservative reanalysis of the same FORS1 data by Wade et al. (2007) did not confirm the longitudinal magnetic field detection at 3σ confidence. Wade et al. (2007) reported longitudinal fields in HD 72106A of 166 ± 70 G from Balmer lines and -11 ± 91 G from metallic lines, and fields in HD 72106B of 52 ± 90 G from Balmer lines and 3 ± 122 G from metallic lines. However, they stressed that ESPaDOnS observations show a clear circular polarization signal, implying that HD 72106A is definitely magnetic.

Due to its young age, probable binarity and magnetic properties, the HD 72106 system is a compelling target for further study. While other binary systems containing a magnetic intermediate-mass star are known, the apparently very young age and wide visual separation of HD 72106 make this system uniquely interesting. We therefore obtained 20 high-resolution spectropolarimetric (Stokes *I* and *V*) observations of the system, with the aim of studying the system's binarity, rotational properties, magnetic field, chemical abundances and surface abundance distributions.

3 OBSERVATIONS

Observations used in this paper were obtained with the ESPaDOnS instrument on the 3.6-m CFHT. ESPaDOnS¹ is a high-resolution echelle spectropolarimeter, with a resolving power of about 65 000 in spectropolarimetric mode, and nearly continuous wavelength coverage from 370 to 1050 nm. This instrument uses a fibre fed design, with a polarimeter module mounted at the Cassegrain focus of the CFHT and an echelle spectrograph located in the Coudé room.

20 Stokes *I* and *V* spectra were obtained over a period of 2 yr, and are summarized in Table 1. 'Normal' readout in spectropolarimetric mode was used, and the atmospheric dispersion corrector was employed for all observations. Magnetic and non-magnetic standard stars were observed as part of this campaign, and produced results in excellent agreement with the literature.

With careful guiding and monitoring, observations of the individual components of HD 72106 were acquired on a few nights with particularly good seeing. This was challenging, however, as the double-star system has a separation of only 0.805 arcsec and ESPaDOnS has a pinhole with a diameter of 1.6 arcsec. On most nights the atmospheric conditions were not good enough for us to resolve the individual components of the double-star system. In these cases the photocentre of the system was observed, providing combined spectra of the system.

The observations were reduced with Libre-ESpRIT (Donati et al. 1997, and in preparation). This is a nearly automatic dedicated data reduction package for ESPaDOnS, which performs a complete calibration and optimal spectrum extraction. Continuum normalization of the resulting unnormalized spectra was performed manually, first by computing a running average on each order of an observed spectrum, then selecting points in the continuum (i.e. maxima) of the smoothed order, and fitting a polynomial (typically of degree 4 or 5) through the selected points. The original observation was then divided by the polynomial to produce a normalized spectral order, and the process was repeated for each order in the echelle spectrum. Similarly, Stokes *V* spectra were normalized by the Stokes *I* continuum polynomial, producing a *V/I_c* spectrum.

¹ See <http://www.cfht.hawaii.edu/Instruments/Spectroscopy/Espadons/> for details.

Table 1. Log of observations of the HD 72106 system obtained with ESPaDOnS. Integration times refer to the complete set of four subexposures used to produce each Stokes V spectrum (Donati et al. 1997). Peak S/N values are reported for the original reduced spectra with 1.8 km s^{-1} spectral pixels. Magnetic measurements for combined observations refer to results for the reconstructed spectra of the primary. The letter D in the LSD detection column indicates that the Stokes V LSD profile is inconsistent with the null field hypothesis at the 99.999 per cent confidence limit, indicating the presence of a magnetic field. The letter N indicates that the inconsistency was at the 99.9 per cent limit or less, and hence cannot be taken as strong evidence for the detection of a magnetic field. The letters MD indicate a ‘marginal’ detection, with a value between 99.9 and 99.999 per cent, and are considered suggestive of a magnetic field, but not conclusive. (Donati, Semel & Rees 1992; Donati et al. 1997). Note that the observation from 2005 February was first published by Wade et al. (2005).

UT date	HJD (−245 0000)	Component (HD 72106)	Integration time (s)	Peak S/N I	B_z (G)	LSD detection
2005 February 22	3423.9248	A and B	2400	201	228 ± 50	D
2006 January 09	3745.029 67	A and B	2400	219	345 ± 42	D
2006 January 11	3747.020 34	A and B	3200	143	261 ± 76	MD
2006 January 12	3747.996 29	A	1200	149	-13 ± 44	MD
2006 January 12	3748.014 96	B	1200	76	2 ± 168	N
2006 February 11	3777.878 60	A and B	2000	238	231 ± 40	D
2006 February 11	3777.951 49	A and B	2400	253	124 ± 37	D
2006 February 12	3778.861 72	A and B	2400	282	282 ± 34	D
2006 February 13	3779.872 02	A and B	2400	184	170 ± 53	D
2006 February 13	3779.981 27	A and B	2400	128	157 ± 156	N
2007 March 02	4161.772 82	A and B	2400	254	350 ± 36	D
2007 March 02	4161.802 56	A and B	2400	265	297 ± 34	D
2007 March 02	4161.902 82	A and B	2400	297	214 ± 30	D
2007 March 03	4162.857 13	A and B	2400	208	246 ± 44	D
2007 March 04	4163.835 61	A and B	2400	322	202 ± 27	D
2007 March 05	4164.846 50	B	3200	209	-51 ± 55	N
2007 March 05	4164.883 87	A	2400	248	320 ± 23	D
2007 March 05	4164.909 61	A and B	1600	277	374 ± 32	D
2007 March 09	4168.857 91	A and B	2400	283	304 ± 33	D
2007 March 09	4168.909 47	A and B	2400	272	252 ± 34	D

4 SPECTRUM RECONSTRUCTION

4.1 Procedure

The two individual spectra of HD 72106A show strong metallic lines with significant variability. In the two observations of HD 72106B on its own we see no metallic absorption line variability, down to the level of the noise in our observations. However, we do see emission in two lines of HD 72106B: clear, variable, emission is present in $H\alpha$, and a small amount of emission in the $O\text{I } 7771 \text{ \AA}$ triplet, as shown in Fig. 1. Thus we confirm the spectroscopic properties of the secondary leading to its classification as a HAeBe star. The emission in HD 72106B varies detectably from night to night. Stronger variations are observed on time-scales of months to years. No detectable change is observed between observations obtained on the same night. A period analysis of the emission line variations yields no significant periodicity in $H\alpha$.

As a consequence of the ~ 0.8 arcsec angular separation between the components of HD 72106, the majority of our spectra of HD 72106A and B are of the combined light from the system rather than light from the individual components. The analysis of observed spectra is much more tractable when one is dealing with light from only one object. Thus we attempted to reconstruct the spectrum of one star from the combined spectra of the system.

As the spectrum of the secondary does not appear to vary (outside of emission lines), we can subtract it from the combined spectra and reconstruct the spectrum of the variable primary. We begin by modelling the normalized flux $I_{\lambda,T}$ at any point in the combined spectrum according to

$$I_{\lambda,T} L_T = I_{\lambda,1} L_1 + I_{\lambda,2} L_2, \quad (1)$$

where $I_{\lambda,1}$ and $I_{\lambda,2}$ refer to the normalized flux spectrum of the primary and secondary, respectively, at wavelength λ , and L_1 and L_2 refer to the luminosity of the primary and secondary, respectively. Thus the product $I_{\lambda,1} L_1$ gives the total flux of the primary. $I_{\lambda,T}$ and L_T are the observed normalized flux spectrum of the system and the total luminosity of the system, respectively. The total luminosity can be written as $L_T = L_1 + L_2$. It is elementary to rearrange equation (1) to obtain the normalized spectrum of the primary, $I_{\lambda,1}$:

$$I_{\lambda,1} = I_{\lambda,T} \left(1 + \frac{L_2}{L_1} \right) - I_{\lambda,2} \frac{L_2}{L_1}. \quad (2)$$

Thus the normalized spectrum of the primary can be reconstructed from an observed spectrum of the system ($I_{\lambda,T}$) and a spectrum of the secondary on its own ($I_{\lambda,2}$). Formally the luminosity ratio should also depend on wavelength; however, since the two components differ in temperature by only 2000 K (see Section 5.2, or Wade et al. 2007) the wavelength dependence is not strong, and thus has been neglected. Nevertheless, the magnitudes used to calculate the luminosity ratio must approximately correspond to the wavelength range of interest in the observed spectrum. The magnitude difference between components (0.62 ± 0.02 in the Tycho V band) is well known from *Hipparcos* and Tycho data (ESA 1997; Fabricius & Makarov 2000), thus the ratio of luminosities at $\sim 5000 \text{ \AA}$ is well established.

4.2 Quality control

To investigate the importance of the assumption of a wavelength-independent luminosity ratio, we used ATLAS9 flux distributions corresponding to the atmospheric parameters of the primary and secondary stars (Table 2). The synthetic flux distributions show that

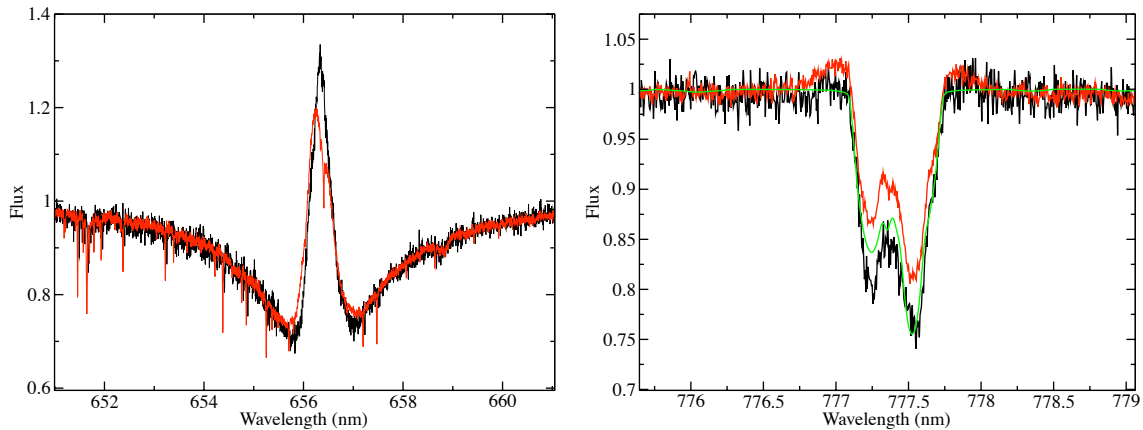


Figure 1. Emission and variability in the H α Balmer line and the O I 7771 Å triplet of HD 72106B (the HAeBe secondary), shown for the two single spectra of HD 72106B. The smooth line in the O I 7771 Å panel is an LTE synthetic profile calculated assuming solar abundances. Strong, variable emission can be seen in H α , while mostly we see only emission infilling in O I 7771 Å. No other lines in the spectrum of the secondary show any evidence of variability.

Table 2. Fundamental physical properties of HD 72106A and B. Magnetic field parameters are determined in Section 8.

	HD 72106A	HD 72106B
T_{eff} (K)	$11\,000 \pm 1000$	8750 ± 500
$\log g$ (CGS)	4.0 ± 0.5	4.0 ± 0.5
L (L_{\odot})	22^{+20}_{-10}	$9.2^{+8.2}_{-4.3}$
R (R_{\odot})	1.3 ± 0.5	1.3 ± 0.5
M (M_{\odot})	2.4 ± 0.3	1.9 ± 0.2
Age (Myr)	6–13	6–13
$v \sin i$ (km s $^{-1}$)	41.0 ± 0.3	53.9 ± 1.0
P_{rot} (d)	0.63995 ± 0.00009	–
i (°)	24 ± 10	–
β (°)	57 ± 5	–
B_p (G)	1230 ± 80	–

the monochromatic luminosity ratio varies by 15 per cent over the ESPaDOnS spectral range. However, since the reconstructed spectrum is essentially a difference spectrum, variations in the luminosity ratio correspond to variations in line depth. Thus a typical line with a depth of 10 per cent of the continuum, at the red or blue end of the spectrum, would have an inaccuracy of ~ 1.5 per cent of the continuum, which is near the noise level in the observations. For most of the lines used in this study this systematic error is even smaller. Comparisons were performed using subtracted spectra calculated assuming different luminosity ratios, as well as comparisons of our reconstructed spectra to observations of the individual components. These showed that systematic errors due to our assumption of a wavelength independent luminosity ratio generally are well below the noise level in our spectra (with the exception of the Balmer lines, which were not included in our analyses).

Stokes V spectra of the primary were reconstructed by attributing the observed polarization signature to the primary and scaling the spectrum based the luminosity ratio of the stars. The assumption that the polarization signal is due solely to the primary is quite reasonable as we see no sign of circular polarization in our spectra of just the secondary. For the purposes of the subtraction, the higher signal-to-noise ratio (S/N) observation of the secondary (HJD 245 4164.846 50) was used.

Contamination of the individual spectra of the secondary by stray light from the primary is potentially a serious source of error in our analysis. However, in our observations this contamination appears to be minimal. Observations of the stars obtained on two different nights, under different seeing conditions, produce identical line profiles (in stable regions of the spectrum) at the level of the noise. If there were a significant amount of contamination, one would expect the contamination to vary with the seeing conditions. Thus the fact that we do not see any such variation suggests that the spectra are not detectably contaminated. A theoretical estimate of the fractional contamination in our spectra, based on the seeing and geometry of the observations, provides further support to this claim. The stars were observed with the centre of the target star on the edge of the pinhole (0.8 arcsec from the pinhole centre) and the other star placed outside the pinhole, ~ 1.6 arcsec away from the pinhole centre. Assuming a Gaussian point spread function for both stars, with a full width at half-maximum equal to the seeing, one can integrate through the pinhole to estimate the flux from each star in our observation. Normalizing the Gaussians by the relative luminosity of the stars, and using the worst seeing conditions for our observations (0.9 arcsec for the primary and 0.7 arcsec for the secondary), we obtain a theoretical worst case contamination of 6 per cent (of the total observed flux) for the observation of the primary and 7 per cent (of the total observed flux) for the observation of the secondary. For contamination by a typical strong line (10 per cent of the continuum), a deviation of roughly 0.6 per cent of the continuum would be seen in the observation, which is at about the noise level of our spectra. Thus we conclude that such contamination is not a serious problem.

Calculation of the spectrum of the primary from equation (2) was performed for all our combined observations of the system, allowing us to reconstruct a sequence of 18 spectra of the primary, typical with a peak S/N ratio of 140 (at 5200 Å). Comparisons of the reconstructed spectra to observations of the primary on its own show excellent concordance in the Balmer lines, and in weakly variable metallic lines, as shown in Fig. 2. Furthermore, this procedure can be applied to recover the spectrum of the secondary (with a typical peak S/N of 90) if we use regions containing only weakly variable lines of the primary. Such experiments also yield good agreement with observations of the secondary on its own. Thus we conclude

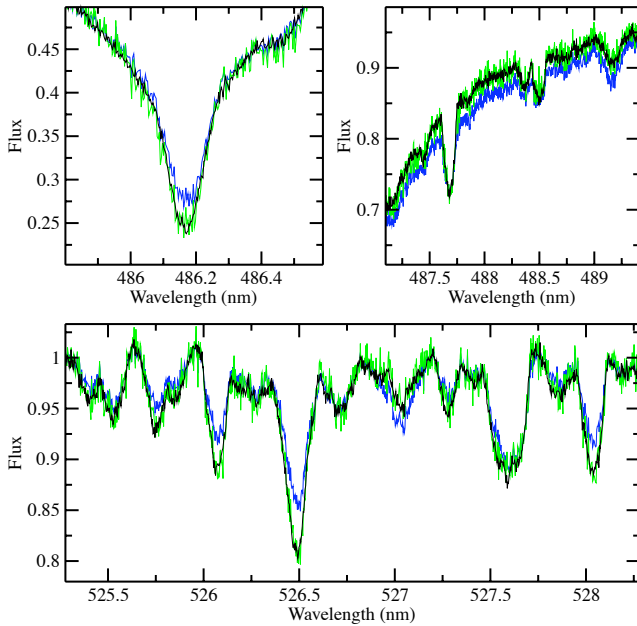


Figure 2. Comparison of an observation of just HD 72106A (black, 2007 March 5), a combined observation of HD 72106A and B (dark grey/blue, 2006 January 11), and the corresponding reconstructed spectrum of HD 72106A (light grey/green) obtained using the procedure described in Section 4. A good match of the reconstructed spectrum to the observation of the primary can be seen in the core (upper left-hand panel) and wings (upper right-hand panel) of $H\beta$, as well as in weakly variable metallic lines (lower panel).

that the spectrum reconstruction is robust, and yields high-precision spectra of the primary component.

5 BINARITY AND EVOLUTIONARY STATE

5.1 Binarity

Determining whether the HD 72106 system is truly a binary or just an optical double star (i.e. an accidental conjunction of two stars at different distances along the line of sight) is critical. If the system is a coeval binary, this fact allows us to constrain the age of the primary much more accurately than would otherwise be possible. Additionally, it suggests that both components formed from approximately the same material, making HD 72106 an interesting system from the point of view of stellar magnetic and chemical evolution.

The system has a projected separation of 0.805 arcsec, and is thus fairly widely separated. Given the *Hipparcos* parallax of 3.60 ± 1.14 mas, this implies a minimum physical separation of 224 ± 71 au.

The *Hipparcos* observations were solved as a binary system producing a ‘good-quality’ solution for the system. This implies that the stars have the same parallax and proper motions, at the precision of the *Hipparcos* observations. *Hipparcos* finds a large proper motion in right ascension of -6.41 ± 0.82 mas yr $^{-1}$ and in declination of 7.97 ± 1.23 mas yr $^{-1}$ (van Leeuwen 2007a). In the spectrum fitting procedure, described in Section 7, we included heliocentric radial velocities for each star as a free parameter. From this fitting, we found identical radial velocities of 22 ± 1 km s $^{-1}$ for both components. Thus it appears that the stars are at the same point in space and moving together in three dimensions, strongly suggesting that they are in fact physically associated.

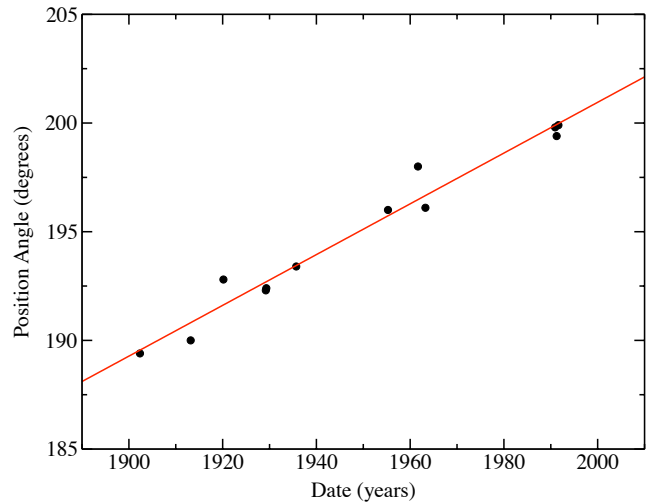


Figure 3. Position angle measurements of the HD 72106, extracted from the WDS, together with their best-fitting line.

Dr Brian Mason at the United States Naval Observatory kindly provided a record of separation and position angle observations from the Washington Double Star Catalogue (WDS), dating back to 1902 (Mason, private correspondence). These have been compiled from a number of sources in the literature, some of which are no longer readily accessible. These observations indicate no significant change in the separation of the components in 89.37 yr, but show a clear and systematic increase in the position angle of the stars of $\sim 10^{\circ}5$, illustrated in Fig. 3. This relative motion was not noted by Wade et al. (2007), presumably because they examined only a subset of these observations. The change in position angle occurs at a rate of $0^{\circ}.117 \pm 0^{\circ}.007$ yr $^{-1}$. In contrast, there is no trend in the set of separation measurements, with an average separation of 0.786 arcsec and a standard deviation of 0.074 arcsec. This kinematical behaviour, together with the identical radial velocities and consistent parallaxes of the components, makes it highly likely that the system is a true binary.

For the system to be a true binary it must be gravitationally bound. We have therefore verified that the observed properties are consistent with binarity. We assumed circular orbits, that the minimum possible separation is the true separation (224 ± 71 au), and used the evolutionary masses derived in Section 5.2. These assumptions are necessary, as we cannot determine the true separation and orbit of the binary without much more accurate observations over a much longer period of time. In this geometry, the relative velocity of the stars is 4.1 ± 0.7 km s $^{-1}$ and the orbital period is 1600 ± 600 yr. This relative velocity is larger than observed, implying that the orbital plane is not parallel to the line of sight, or that the stars’ physical separation is larger than we have assumed (which it almost certainly is). If the orbital plane of the stars is in the plane of the sky, the orbital period implies that an observed change in position angle of $0^{\circ}.219 \pm 0^{\circ}.079$ yr $^{-1}$ should be seen, which is roughly consistent with the WDS observed rate of change. Thus we conclude that a wide range of orbits are consistent with the observed radial velocities and position angle change. One example of an orbit that is consistent with all observations is a circular orbit with a semimajor axis of 350 au (implying a period of 3200 yr), with an inclination of the orbital plane to the line of sight of $44^{\circ}5$, and a position angle of the orbital axis of 200° . In this scenario, the two components would both be crossing the line of sight at the current epoch, yielding

zero relative radial velocity. The change in position angle would be 0.112yr^{-1} and a apparent separation would be 0.807 arcsec. In this example, in 800 yr the components would achieve their maximum difference in relative radial velocity of 2.3 km s^{-1} , and their maximum apparent separation of 1.151 arcsec. Of course, there is no reason to prefer this particular solution over the many others that are consistent with observations.

When placed on the Hertzsprung–Russell (HR) diagram, as discussed in Section 5.2, the stars are found to have positions that allow for a range of coeval solutions. Thus according to the HR diagram positions, it is quite possible for the stars to have formed and evolved together, further supporting the binary hypothesis.

Thus we conclude that HD 72106 is very likely a true binary system. The stars are at the same position in space, moving in the same (three-dimensional) direction at the same velocity, allow for a wide range of gravitationally bound orbits, and appear to be consistent with coeval evolution.

5.2 Fundamental physical properties

Effective temperature and surface gravity for both components of HD 72106 were determined by fitting Balmer lines. Unfortunately, no intermediate-band photometry of the components of HD 72106 was available. The $H\alpha$ Balmer line in HD 72106B is clearly contaminated with emission, and hence was not considered in this analysis. The other Balmer lines of both stars appear to be free of emission, and lines of each star can be well fitted with a single model atmosphere.

Wade et al. (2005) performed Balmer line fitting for both components of HD 72106 using observations from the FORS1. For the primary they found: $T_{\text{eff}} = 11\,000 \pm 1000\text{ K}$ and $3.5 \leq \log g \leq 4.5$ (best-fitting $T_{\text{eff}} = 11\,000\text{ K}$ $\log g = 4.0$). For the secondary they found $T_{\text{eff}} = 8000 \pm 500$ and $4.0 \leq \log g \leq 4.5$ (best-fitting $T_{\text{eff}} = 8000\text{ K}$, $\log g = 4.5$). We repeated the fitting procedure with the FORS1 spectra, using the method outlined above. Solar abundance ATLAS9 model atmospheres were used (Kurucz 1993) to produce the synthetic Balmer lines and the model lines were convolved with a Gaussian instrumental profile of appropriate width to match the observations. We arrived at results identical to Wade et al. (2005) for the primary. For the secondary we find the temperature range $7500\text{--}9000\text{ K}$ and the $\log g$ range $4.0\text{--}4.5$ more realistic, with a best-fitting value of 8000 K at $\log g = 4.5$. These values from the FORS1 observations are also consistent with those determined from Balmer line fitting of our ESPaDOnS spectra.

In Section 7 we performed detailed spectrum synthesis of both components of HD 72106 and found that the atmospheric parameters of the secondary derived from Balmer lines are not compatible with the metallic line spectrum. In particular, we were unable to simultaneously fit lines of Fe I and Fe II, Cr I and Cr II, and Ti I and Ti II, suggesting a problem with the local thermodynamic equilibrium (LTE) ionization balance of these species. Including T_{eff} and $\log g$ as free parameters in the spectrum fit allowed us to satisfactorily match the metallic line spectrum, provided that $T_{\text{eff}} \sim 8750\text{ K}$ and $\log g \sim 4.0$. While these values produce a somewhat poorer fit to the Balmer lines, the fit to metallic lines is improved substantially. Hence we conclude that the higher temperature and lower $\log g$ provide a better description of the atmosphere of HD 72106B.

Including the full range of T_{eff} and $\log g$ which provide acceptable fits to both the metallic line spectra and the Balmer lines of HD 72106B, we adopt $T_{\text{eff}} = 8750 \pm 500\text{ K}$ and $\log g = 4.0 \pm 0.5$.

Using the *Hipparcos* parallax of the HD 72106 system and the Tycho magnitudes (Fabricius & Makarov 2000) we determined the

luminosities of both components of HD 72106. To convert the observed Tycho V magnitudes ($V_T = 9.00 \pm 0.01$ for the primary and $V_T = 9.62 \pm 0.02$ for the secondary) into a Johnson V magnitudes (V) we used the empirical relation (ESA 1997)

$$V = V_T + 0.09(B_T - V_T), \quad (3)$$

where B_T is the Tycho B magnitude. For the secondary, the bolometric correction for main-sequence stars from Gray (2005) was used, yielding a value of $0.01 \pm 0.06\text{ mag}$. For the primary, the bolometric correction relation of Landstreet et al. (2007) was used, yielding a value of $-0.37 \pm 1.9\text{ mag}$. This calibration is tailored specifically for magnetic chemically peculiar A and B type (Ap and Bp) stars. As will be shown, the primary has properties very similar to those of a Bp star, thus the calibration of Landstreet et al. (2007) is more appropriate. From this procedure we find that the luminosity of the primary is $22^{+20}_{-10} L_\odot$, and the luminosity of the secondary is $9.2^{+8.2}_{-4.3} L_\odot$. From the Stefan–Boltzmann equation we find identical radii, given our uncertainty, of $1.3 \pm 0.5 R_\odot$ for the primary and $1.3 \pm 0.5 R_\odot$ for the secondary (with 1σ error bars).

With the luminosity and effective temperature determined, we can place the stars on the HR diagram, as shown in Fig. 4, and compare their positions with theoretical evolutionary tracks and isochrones. Using the pre-main-sequence evolutionary model calculations from CESAM (version 2K; Morel 1997) and the birth line of Palla & Stahler (1993), we find the mass of the primary to be $2.4 \pm 0.3 M_\odot$ and mass of the secondary to be $1.9 \pm 0.2 M_\odot$ (with 1σ error bars). From this we derive an evolutionary $\log g$ of 4.6 ± 0.3 for the primary, and 4.5 ± 0.3 for the secondary, both of which are consistent with (although systematically higher than) the spectroscopic values. We determine the binary’s age to be between 6 and 13 Myr (measured from the birth line), based on the position of the secondary in the HR diagram and assuming a protostellar accretion rate of $10^{-5} M_\odot \text{ yr}^{-1}$. In doing this we have assumed, as most investigators do, that the presence of circumstellar material implies that the H A e Be secondary is on the pre-main sequence, and thus that we can constrain its age with pre-main-sequence evolutionary tracks. Given the derived admissible range of age and mass, the primary could still be on the pre-main sequence, reaching the zero-age main-sequence (ZAMS) in less than 1 Myr. The

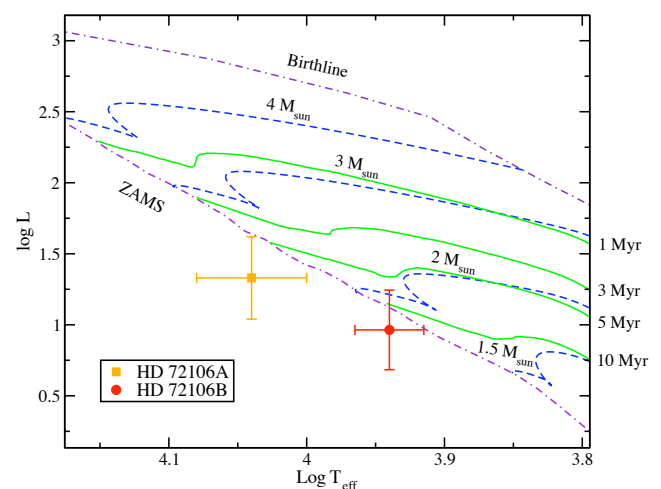


Figure 4. An HR diagram containing both components of the HD 72106 system. Isochrones (solid lines), evolutionary tracks (dashed lines) and the ZAMS line, are calculated from CESAM models (Morel 1997). The birth line (for an accretion rate of $10^{-5} M_\odot \text{ yr}^{-1}$) is taken from Palla & Stahler (1993).

most likely case (from ‘best-fitting’ positions) is that the system is ~ 10 -Myr old, and the primary has just entered the main sequence, while the secondary is on its final approach to the ZAMS. In this case the primary would have spent ~ 6 Myr on the main sequence, giving it a fractional age on the main sequence (τ) of 0.01. In the oldest limiting case, the secondary is just reaching the ZAMS and the primary has been on the main sequence for ~ 9 Myr giving it a fractional age of 0.015. Further observations, particularly a more accurate distance measurement, are necessary to more precisely determine the evolutionary status of the primary. Regardless, while HD 72106A may not be on the pre-main sequence, it is certainly one of the youngest known main-sequence stars of its type.

The physical properties we derive for both stars are summarized in Table 2. It is worth noting that, while the absolute luminosities of the components are poorly determined, their ratio ($L_A/L_B = 2.3 \pm 0.4$) is very well determined. This is because the major uncertainty in the absolute luminosity is the distance to the stars, and the stars are located at the same distance. Thus the values of the radii or masses derived for the two stars are not independent, since spacing between the components in $\log L$ on the HR diagram must remain fixed.

6 LEAST-SQUARES DECONVOLUTION AND STELLAR ROTATION

6.1 LSD and longitudinal magnetic field

In order to measure the magnetic fields of HD 72106A and B, we employed least-squares deconvolution (LSD; Donati et al. 1997). This cross correlation technique uses a table of input atomic data (the line mask) to produce a deconvolved mean line profile from an observed spectrum. Using a large number of lines, ~ 7000 for our spectra, can result in dramatic improvements in S/N. LSD was performed on all our observations of the individual components of HD 72106, as well as the reconstructed spectra of HD 72106A. The full set of LSD profiles in Stokes I and V for HD 72106A, calculated using specific elements, are presented in Fig. 10 in Section 8.

The line mask used for the primary was calculated assuming an effective temperature of 11 000 K, $\log g$ of 4.0, typical Ap star abundances, and a line depth cut-off of 0.1 (as a fraction of the continuum). This produced LSD profiles with a mean wavelength of 523.4 nm, a mean Landé factor of 1.27, and a mean excitation potential of 5.97 eV, with small variations between observations. These mean values represent the average values of the lines used in the computation of the LSD profile, weighted by the S/N of the lines. The computed LSD profiles are relatively insensitive to variations or errors in the line mask in temperature or abundance (e.g. Shorlin et al. 2002). For the secondary, a $T_{\text{eff}} = 8750$ K, $\log g = 4.0$ line mask was used with solar abundances and a line depth cut-off of 0.1. With this mask we found a mean wavelength of 533 nm, a mean Landé factor of 1.21, and a mean excitation potential of 3.99 eV.

The longitudinal magnetic field was then measured from each set of LSD profiles, using equation (1) of Wade et al. (2000). Integration was performed through the portion of the line profile that exceeded 15 per cent of the total line depth, with an additional 5 km s^{-1} on either end. The longitudinal field data are presented in Table 1. We find a maximum longitudinal field for HD 72106A of 374 ± 32 G.

In the secondary, in contrast, we detect no longitudinal field, with error bars of 50–150 G. Moreover, the LSD profiles of HD 72106B yield no evidence of a magnetic field. The analysis of HD 72106B by Wade et al. (2007), discussed in Section 2, found no magnetic

field in either the Balmer lines or metallic lines with error bars of ~ 100 G. Thus we conclude that if the secondary has a magnetic field, its longitudinal component never exceeds ~ 200 G.

6.2 Rotation period of HD 72106A

Variability in the spectrum of HD 72106A was noted by Wade et al. (2005). A careful re-examination of their data shows that most of their reported spectrum variability was due to the lack of atmospheric dispersion correction during the acquisition of one of their spectra. However, we do see substantial metal line variability in all of our spectra of the primary, with a variation time-scale on the order of a day. We do not confirm the rotational period of approximately 2 d proposed by Wade et al. (2005).

In order to investigate whether this variability in HD 72106A could be due to rotational modulation, we fit the longitudinal magnetic field observations with a first order sinusoid of fixed period, but free phase, amplitude and mean. This process was repeated for many periods to construct a periodogram in a fashion analogous to the Lomb–Scargle method (Press et al. 1992). In the resulting periodogram, shown in Fig. 6 (left-hand frame), we found a number of periods that provided a good fit to our observations, with the best period at about 0.640 d. Our longitudinal field observations, phased with this best-fitting period, are shown in Fig. 5. We interpret this variability as being due to the rotation of HD 72106A, in the context of the oblique rotator model (ORM; e.g. Stibbs 1950). Using the radius determined in Section 5.2, together with the best-fitting period, we find that a range of inclination angles will satisfy the derived $v \sin i$ from Section 7.1. The smooth sinusoidal behaviour of the longitudinal magnetic field, the stability of the period over two years, and the period consistent with the star’s $v \sin i$ and radius all support our interpretation of the variability in the context of the ORM. It is worth noting that the phase coverage for periods around 0.64 d is poor between phases 0.4 and 0.7. However, the remainder of our analysis is based on rotationally broadened LSD profiles rather than longitudinal field measurements, which mitigates this phase gap.

In order to search further for periodic variations in the observations of HD 72106A, we examined the variability within the Stokes I and V LSD profiles. To do this, we constructed a periodogram

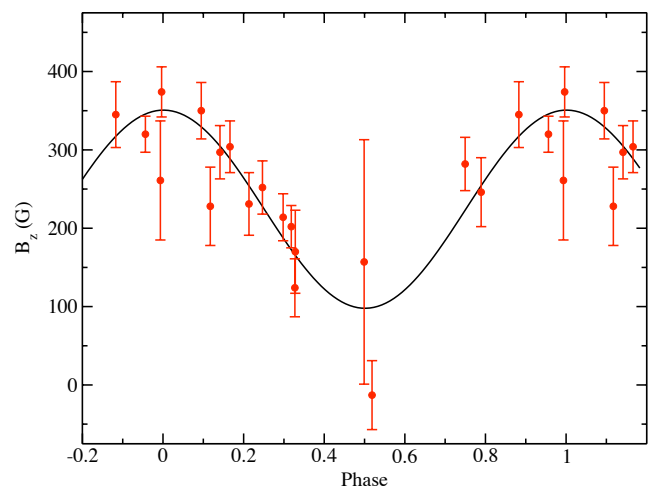


Figure 5. Longitudinal magnetic field measurements of HD 72106A (points), phased with the adopted 0.63995 d period, and the longitudinal field curve from the best-fitting dipolar magnetic geometry (solid line) determined in Section 8.

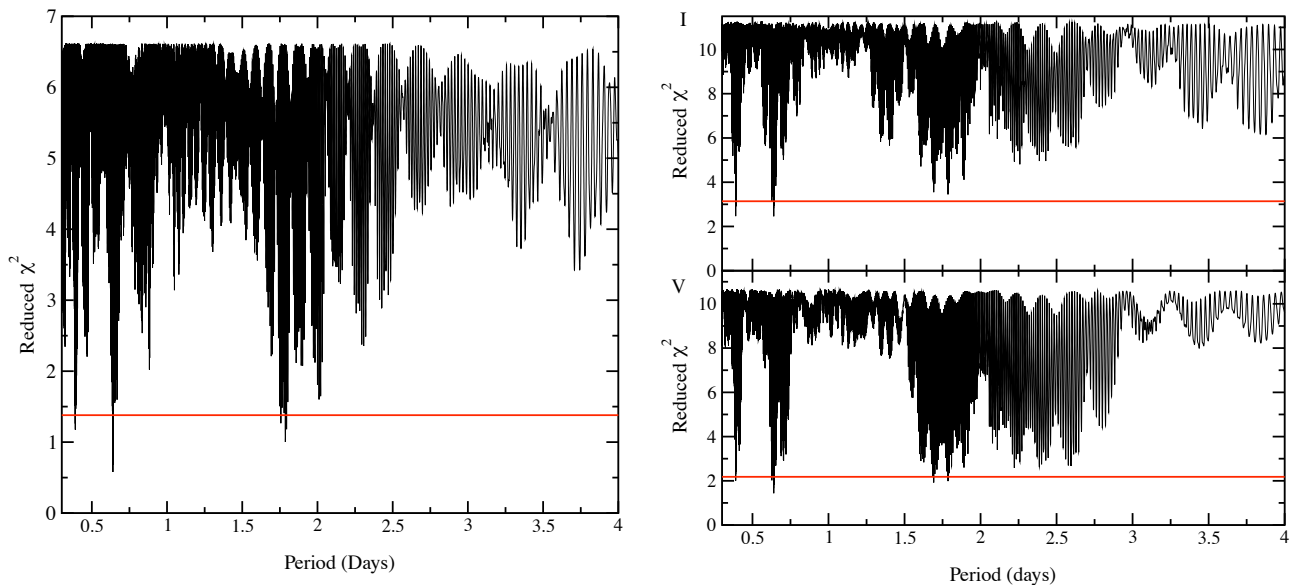


Figure 6. Periodograms for HD 72106A based on longitudinal field measurements (left-hand panel) and LSD profile variability (right-hand panel) in Stokes *I* and *V*. A few minima in each periodogram fall below the 99 per cent confidence limit (horizontal line); however, the deepest minimum is at 0.64 d in all cases.

based on the variability of each pixel of the LSD profiles, using the same technique as for the longitudinal field measurements. The periodograms for each pixel were then averaged and weighted by the amplitude of variability in that pixel, to produce a periodogram for the entire set of LSD profiles. This procedure was performed for both Stokes *I* and Stokes *V* LSD profiles, and the resulting periodograms are shown in Fig. 6 (right-hand frame).

Periods at 0.389 83, 0.639 95, 1.6921 and 1.7859 d all produced minima in one or more of the periodograms that were within the 99 per cent confidence limit of the global minimum. These candidate periods were each investigated by eye. Each period was used to determine the phasing of the LSD profiles, and the profiles were examined for variability physically consistent with rotation. For a physically sensible rotational variation, features must move smoothly across the profile from blue to red, and there must not be large changes in the profile with small changes in phase. Upon careful examination, the period of 0.639 95 d was the only period to produce a physically sensible phasing of the *I* and *V* profiles. Thus we adopt the rotational ephemeris (with maximum longitudinal field at zero phase) for HD 72106A:

$$\text{HJD} = (245\,3747.017 \pm 0.013) + (0.639\,95 \pm 0.000\,09)E. \quad (4)$$

The LSD profiles phased according to this ephemeris are shown in Fig. 10. This period is notable in that it is one of the shortest rotation periods seen in any magnetic intermediate-mass star.

7 SPECTRUM SYNTHESIS

The ZEEMAN2 spectrum synthesis code was used to model the observed spectra of HD 72106A and B. This code solves the polarized radiative transfer equations in LTE. In order to find the best-fitting model of the observed spectrum, a Levenberg–Marquardt χ^2 minimization routine was implemented. This routine consistently produced high-quality, stable solutions with good efficiency. The results were, in all cases, carefully examined by eye to verify that they represented both a physical solution and the true best fit, rather than a local minimum in χ^2 .

Chemical abundances [in the form (N_X/N_{tot}) , assuming $N_{\text{He}}/N_{\text{H}} = 0.098$], projected rotational velocity ($v \sin i$), and for HD 72106B, microturbulence, were the free parameters in the fit. Chemical abundances were assumed to be uniform both vertically and horizontally. This is not strictly true for HD 72106A, which displays patchy distributions of some elements, discussed further in Section 8. Thus the abundances derived for HD 72106A are averages across the visible disk of the star at the phase of the observation used (the 2006 January 12 observation, rotation phase 0.53). However, since there is not much change in equivalent widths between spectra, the observation of HD 72106A used is approximately representative of the global average for all elements, even those with patchy distributions.

Microturbulence is not expected in the primary, due to its strong magnetic field. By analogy to Ap/Bp stars, the magnetic field would likely suppress microturbulence, thus microturbulence was fixed at 0 km s^{-1} for HD 72106A. This did not produce any detectable discrepancies between our model and observed spectra. For the primary we adopted the magnetic field model derived from the Stokes *V* LSD profiles in Section 8. Uncertainties in this magnetic field model have a negligible impact on the derived abundances. For the secondary no magnetic field was included.

Standard model atmospheres with solar abundance, computed with the ATLAS9 (Kurucz 1993) code, were used as input for ZEEMAN2. Input atomic line data were obtained from the Vienna Atomic Line Database (VALD) (Kupka et al. 1999), using the ‘extract stellar’ utility in VALD (with the default line list configuration).

Seven segments of spectrum, ranging from 100 to 200 Å in length, were independently fit for each star (4170–4265, 4400–4500, 4500–4600, 4600–4700, 5000–5200, 5200–5400 and 5400–5600 Å). The final best-fitting parameters were taken as the average of the parameters derived from individual windows. Uncertainties were taken as the standard deviation of the parameters. For abundances that were only constrained in a few windows, the uncertainties were estimated by eye, by taking the change in abundance necessary to shift the synthetic spectrum beyond the noise in the observed spectrum and any possible continuum normalization errors. In these

Table 3. Averaged best-fitting chemical abundances, $v \sin i$, and microturbulence (ξ) for HD 72106A and B as well as solar abundances from Grevesse et al. (2005). Entries marked by an asterisk (*) are based on only a few lines.

	HD 72106A	HD 72106B	Solar
$v \sin i$ (km s ⁻¹)	41.0 ± 0.6	53.9 ± 1.0	
ξ (km s ⁻¹)		2.3 ± 0.6	
He	-2.8 ± 0.2*		-1.07
C		-3.40 ± 0.08	-3.61
O	-3.0 ± 0.15*		-3.34
Mg	-4.0 ± 0.3	-4.60 ± 0.16	-4.47
Al	-5.8 ± 0.2*		-5.63
Si	-3.60 ± 0.14	-4.7 ± 0.3	-4.49
Ca	-5.1 ± 0.3*	-6.0 ± 0.2	-5.69
Sc	-8.6 ± 0.2*	-9.13 ± 0.08	-8.83
Ti	-6.04 ± 0.10	-7.24 ± 0.09	-7.10
Cr	-4.33 ± 0.13	-6.32 ± 0.19	-6.36
Mn		-6.7 ± 0.3*	-6.61
Fe	-3.49 ± 0.07	-4.64 ± 0.17	-4.55
Ni		-6.3 ± 0.3	-5.77
Sr	-8.7 ± 0.4*	-9.2 ± 0.4*	-9.08
Y		-10.0 ± 0.3	-9.79
Ba	≤ -8.5*	-10.2 ± 0.4*	-9.83
Ce		≤ -9.0*	-10.30
Nd	-7.5 ± 0.3*	≤ -9.2*	-10.54

case the presented uncertainties are indicated by an asterisk (*) in Table 3.

7.1 Abundances in HD 72106A

Sample fits of our synthetic spectra to the observations of HD 72106A are presented in Fig. 7. The average best-fitting abundances and $v \sin i$, with uncertainties, are given in Table 3. The abundances are presented graphically, relative to solar abundances, in Fig. 8.

Remarkably strong overabundances of Si, Cr, Fe and Nd are found in HD 72106A. The Si and Fe abundances appear to be above solar by ~ 1 dex, and Cr appears to be enhanced by ~ 2 dex, whereas He appears to be ~ 1.5 dex underabundant. A number of elements,

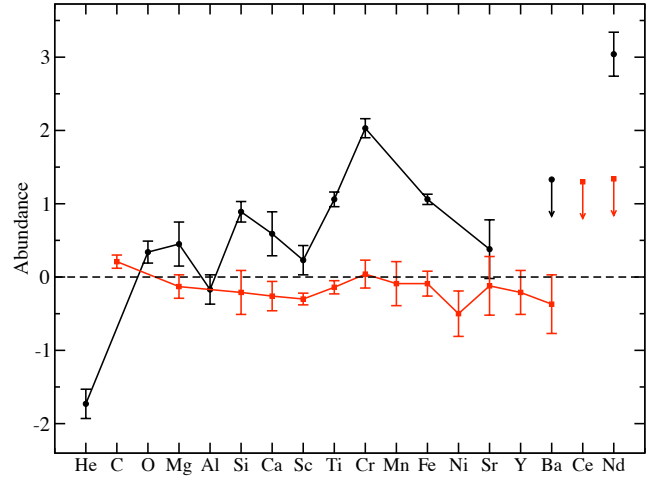


Figure 8. Abundances relative to solar for HD 72106A (circles) and HD 72106B (squares), averaged over all spectral windows modelled. The dashed line at 0 represents solar abundance, based on Grevesse, Asplund & Sauval (2005). Points marked with an arrow indicate the value is an upper limit only. Strong departures from solar abundance can be seen for HD 72106A, whereas HD 72106B has largely solar abundances.

such as Al, Sc and Sr appear to have solar abundances. A couple of elements, particularly Mg and O, hint at possible peculiarity but require further study before concrete conclusions can be drawn. We find $v \sin i = 41.0 \pm 0.7$ km s⁻¹, which is fairly low for a main-sequence B star, but within the normal range for Bp stars. The strong overabundances in Si, Cr, Fe and Nd, as well as the underabundance in He, are common features of cooler Bp stars (Jaschek & Jaschek 1995).

7.2 Abundances in HD 72106B

HD 72106B was initially modelled with an effective temperature of 8000 K and $\log g = 4.5$. However, as discussed in Section 5.2, this model provided a poor fit to the observed stellar spectrum, and hence T_{eff} and $\log g$ were included as free parameters in our fit. Chemical

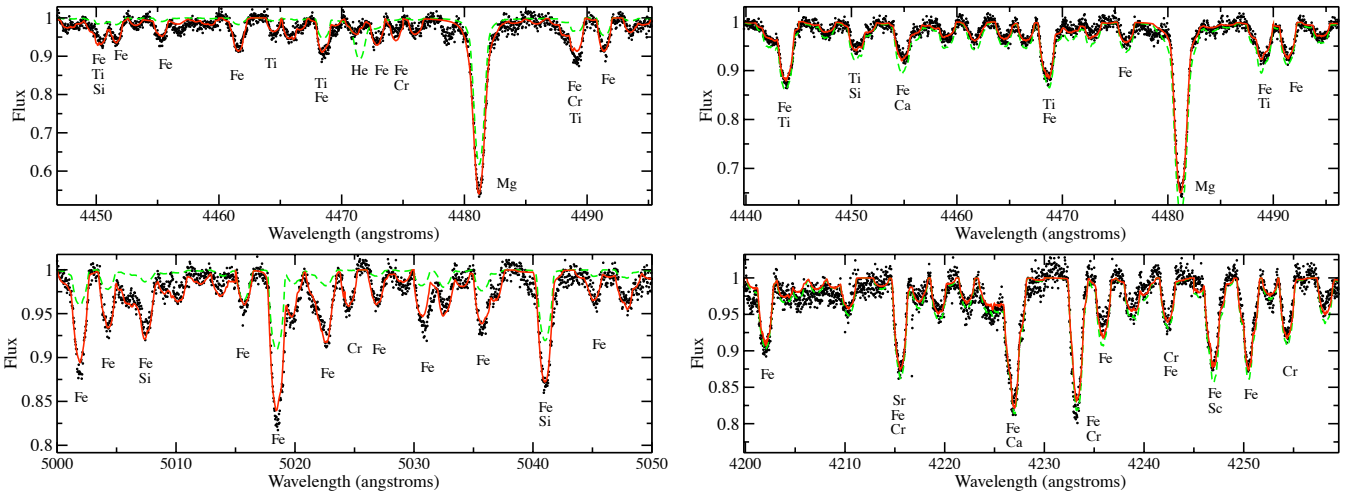


Figure 7. Sample best-fitting synthetic spectra for HD 72106A (left-hand panels) and HD 72106B (right-hand panels) in two independently fit spectral windows. The observations of the individual stellar components from 2007 March 5 are used for both stars. Major contributors to each line have been labelled, in order of importance. The smooth solid line is the best-fitting spectrum in this region, the dashed line is a spectrum computed with solar chemical abundances. Dots represent the observations.

abundances, $v \sin i$ and microturbulence were also included as free parameters in the fit, while the magnetic field was set to zero. From this procedure we found a best-fitting T_{eff} of 8750 ± 500 K and $\log g = 4.0 \pm 0.5$. Sample best-fitting synthetic spectra can be seen in Fig. 7, compared with the observed spectrum used in the fitting process. The average best-fitting abundances, $v \sin i$, and microturbulence derived for HD 72106B are presented in Table 3, and the abundances are shown graphically relative to solar abundances in Fig. 8.

The large majority of elements are consistent with solar abundances, within 2σ at most. A few elements appear to depart marginally from solar values, with a significance slightly greater than 2σ . C appears to be overabundant with $\sim 2\sigma$ significance, whereas Sc is underabundant by $\sim 3\sigma$. Under abundances of Sc and Ca are characteristic of Am stars. However, since the Ca abundance is approximately normal in HD 72106B, and iron peak elements are not overabundant, HD 72106B is not Am star. Thus, nearly all the elements are within 2σ of solar abundance, and no elements display the strong peculiarities seen in HD 72106A. This result is consistent with the approximately solar abundances seen in most HAeBe stars (Acke & Waelkens 2004). It is of possible importance that the best-fitting abundances are consistently below solar, albeit with small significance levels. However, there remains a large uncertainty in our best-fitting temperature and $\log g$. An increase in the adopted temperature by 250–9000 K would systematically increase the best-fitting abundances by 0.1 or 0.2 dex, essentially eliminating this trend.

8 MAGNETIC DOPPLER IMAGING

The observed metallic line variability with rotation strongly suggests that there are horizontal chemical abundance inhomogeneities in the atmosphere of HD 72106A. In an effort to determine the structure of these inhomogeneities, as well as the magnetic field structure, we performed magnetic Doppler imaging (MDI) on HD 72106A. Doppler imaging is a method for inverting a time-series of variable line profiles at known rotational phases in order to reconstruct the surface distributions that give rise to the observed variations. In MDI, polarized line profiles are also included in the inversion, allowing for the reconstruction of a magnetic field through the Zeeman effect.

We used the MDI code *INVERS10*, developed by O. Kochukhov and N. Piskunov (Kochukhov & Piskunov 2002; Piskunov & Kochukhov 2002) for surface chemical abundance mapping and magnetic field reconstruction. This program performs accurate polarized LTE spectrum synthesis, using pre-calculated model stellar atmospheres. *INVERS10* allows for simultaneous modelling of multiple chemical elements and multiple wavelength regions, and takes into account blended lines. Tikhonov regularization is used (Tikhonov 1963; Piskunov & Kochukhov 2002), to constrain abundance gradients in the surface map. While Tikhonov regularization can be applied to the magnetic field as well, we opted for multipolar regularization (Kochukhov et al. 2002), due to the lack of Stokes Q and U spectra and the relatively low S/N in the Stokes V spectra.

Input into the MDI model of HD 72106A included the adopted effective temperature and surface gravity from Section 5.2: $T_{\text{eff}} = 11\,000$ K and $\log g = 4.0$. The inclination angle of the star's rotation axis to the line of sight (i) was calculated using the radius determined in Section 5.2, the rotation period from Section 6.2, and the $v \sin i$ determined from spectrum fitting in Section 7.1, giving an angle of $i = 24^\circ \pm 10^\circ$. The phasing of observations was determined from the adopted ephemeris. An initial $v \sin i$ of 41 km s^{-1} and abun-

dances, for the treatment of blended lines, were taken from the results of our abundance analysis of HD 72106A, presented in Section 7.1. The initial magnetic field geometry was assumed to be a dipole.

Initially, five Si II lines at 4128, 4130, 5056, 5041 and 6371 Å were considered for Doppler imaging. However, the S/N of the observations was insufficient to produce high-quality Doppler maps for Si, as well as for all other elements.

In order to improve the S/N of our observations we calculated LSD profiles for individual chemical elements, and then performed MDI using those profiles. The use of LSD profiles to determine magnetic field geometries, stellar pulsation and surface structures (Donati et al. 2000, 2001) is well established. LSD profiles of individual elements were obtained by creating LSD line masks containing only lines of the element of interest. The line masks used were derived from the 11 000 K Ap star line mask discussed in Section 6. Doppler imaging was then performed using the set of LSD profiles for each element as described above. The elements Si, Ti, Cr and Fe were used in this process. Due to the low S/N in Stokes V for the purposes of MDI, even in the LSD profiles, the magnetic field geometry was determined from the higher S/N Cr and Fe profiles, and held fixed for the lower S/N Si and Ti profiles. Additionally, the multipolar magnetic field regularization was restricted to $l = 1$ modes, effectively providing a 'dipolar regularization'. Mean LSD profile atomic data were used for the Doppler imaging process, providing wavelengths and excitation potentials, shown in Table 4. Since the depth of the LSD profile is a complex function of the lines used in the analysis, the mean gf value may not represent a realistic oscillator strength for the LSD profile. As a consequence, the absolute abundance scale of the Doppler images is somewhat uncertain.

The final abundance maps, and magnetic modulus and vector maps, are shown in Fig. 9. The corresponding best fits to the LSD line profiles are presented in Fig. 10. In this process we have assumed that the LSD profile behaves like a real spectral line, and that the mean atomic data are approximately representative of the LSD profile. Additionally, unaccounted for line blends in the LSD process could slightly distort the LSD profiles, adding some uncertainty to the finer details of the Doppler map. Thus, while the large-scale features of the Doppler imaging process are fairly certain, one should use some caution in interpreting the finer details of the map, as well as the absolute abundance scale. Similarly, the use of LSD profiles to reconstruct the magnetic field geometry through MDI could potentially introduce some systematic uncertainty into the derived geometry. However, the magnetic geometry derived through MDI is fully consistent with the longitudinal magnetic field variability and nearly identical to that we derive by directly modelling the LSD Stokes V profiles with a simple dipole.

Strong inhomogeneities are reconstructed for all four elements. Ti, Cr and Fe all seem to share similar abundance patterns, though Cr appears to have somewhat larger, more elongated spots. The similar distributions are reflected in the similar phase variations of

Table 4. Mean atomic data for the element specific LSD profiles used in MDI.

	Si	Ti	Cr	Fe
Mean wavelength (Å)	5580	4970	5190	5310
Mean excitation (eV)	7.68	2.03	5.65	6.46
log mean gf	−0.17	−0.63	−0.26	−0.35
Mean Landé factor	1.35	1.12	1.29	1.27

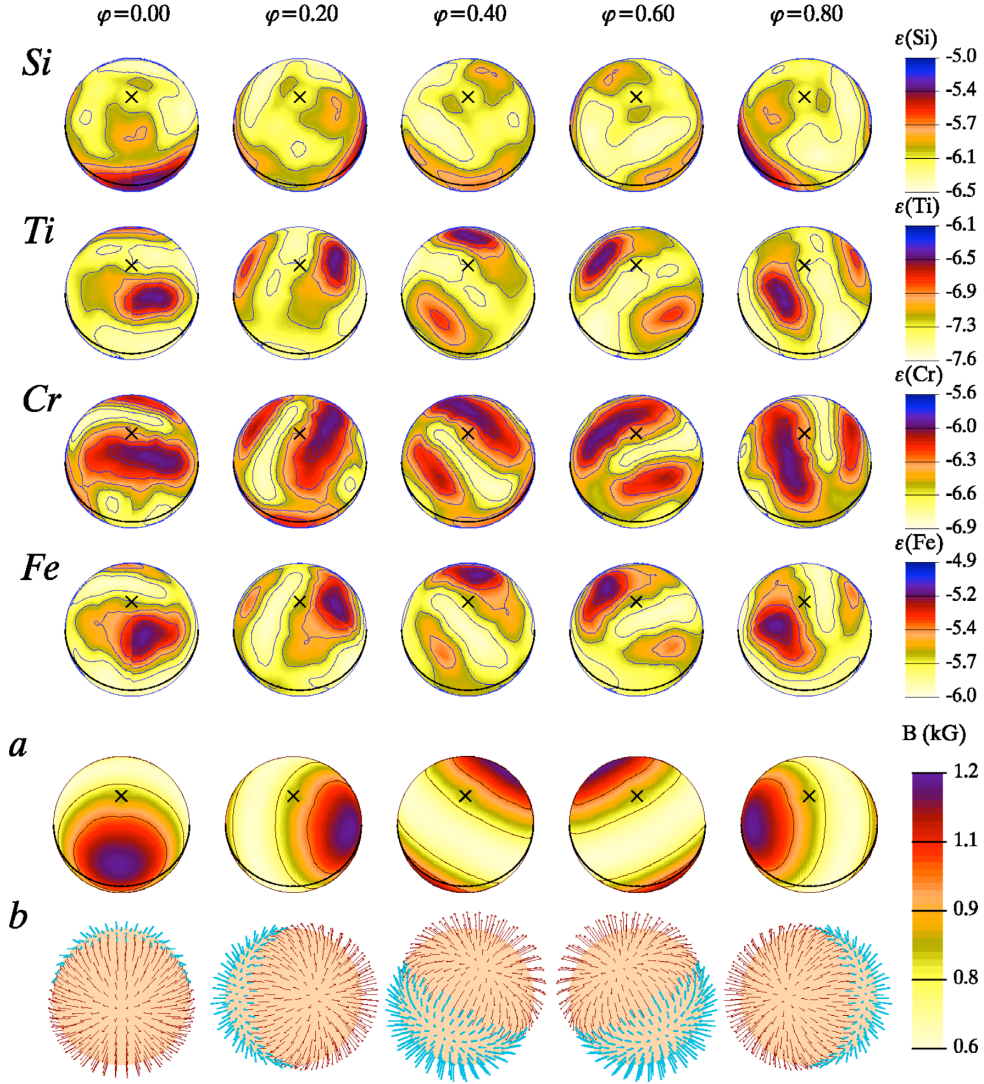


Figure 9. Surface maps of Si, Ti, Cr and Fe abundances, as well as the magnetic field, for HD 72106A. The maps are all based on fits to LSD profiles. The 'X' represents the rotational pole, the thick line circle indicates the rotational equator. The scale on the right-hand side in the abundance maps is in units of $\log(N_X/N_{\text{tot}})$. The map of the magnetic field intensity is labelled 'a' and the magnetic field direction is 'b' (shown as arrows).

the LSD profiles in Fig. 10. A large patch of overabundance centred near phase 0 is apparent in all three maps, with another somewhat smaller overabundance spot about 180° away in longitude, at the same latitude, around phase 0.6. The Si map shares the larger spot but not the smaller, and this is reflected in the Si LSD profiles. There appears to be a large overabundance spot of Si at the equator (seen near the limb) near phase 0, although the sensitivity of the map is poor in that region due to its small projection. Note that the fits to the relatively noisy Si LSD profiles are somewhat poorer than the fits for the other elements analysed, thus the results for Si are somewhat more uncertain. The magnetic field geometry derived, shown in Fig. 9, is essentially described by a dipole with a strength at the magnetic pole (B_p) of 1230 ± 80 G, and an obliquity angle of the magnetic dipole with respect to the rotation axis (β) of $57^\circ \pm 5^\circ$. This dipole magnetic field geometry is in full agreement with the observed longitudinal magnetic field variability. If there are any departures from a purely dipolar field geometry, they are not evident in our data. When the magnetic field geometry is compared to the abundance maps, it appears that the positive magnetic pole

lies near the large spot of overabundance at phase 0 in all four maps. However, the magnetic pole is offset from the centre of the abundance spot, being nearer the rotational equator and at a slightly different latitude, thus the relationship is not entirely clear.

Interestingly, the pattern of abundances seen here for HD 72106A bear some similarities to those obtained for ϵ Ursae Majoris by Lueftinger et al. (2003). ϵ Ursae Majoris (HD 112185) is a 9000 K ($\log g = 3.6$) late main-sequence Ap star with ~ 5 d period and a dipole field strength of several hundred gauss. Lueftinger et al. (2003) constructed Doppler maps of Ti, Cr and Fe, as well as Ca, Mg, Sr and Mn. They found distributions of Cr and Fe very similar to each other, with two large spots of overabundance near the longitude of the magnetic poles. Ti was roughly anticorrelated with Fe and Cr, displaying two large spots of underabundance at the same positions as the overabundance spots of Cr and Fe. We see little variability in Ca, and have insufficient S/N in Ca, Mg, Sr and Mn to construct Doppler maps for comparison with Lueftinger et al. (2003). The results for Fe and Cr are qualitatively very similar to what we see for HD 72106A; however, the results for Ti differ significantly. The

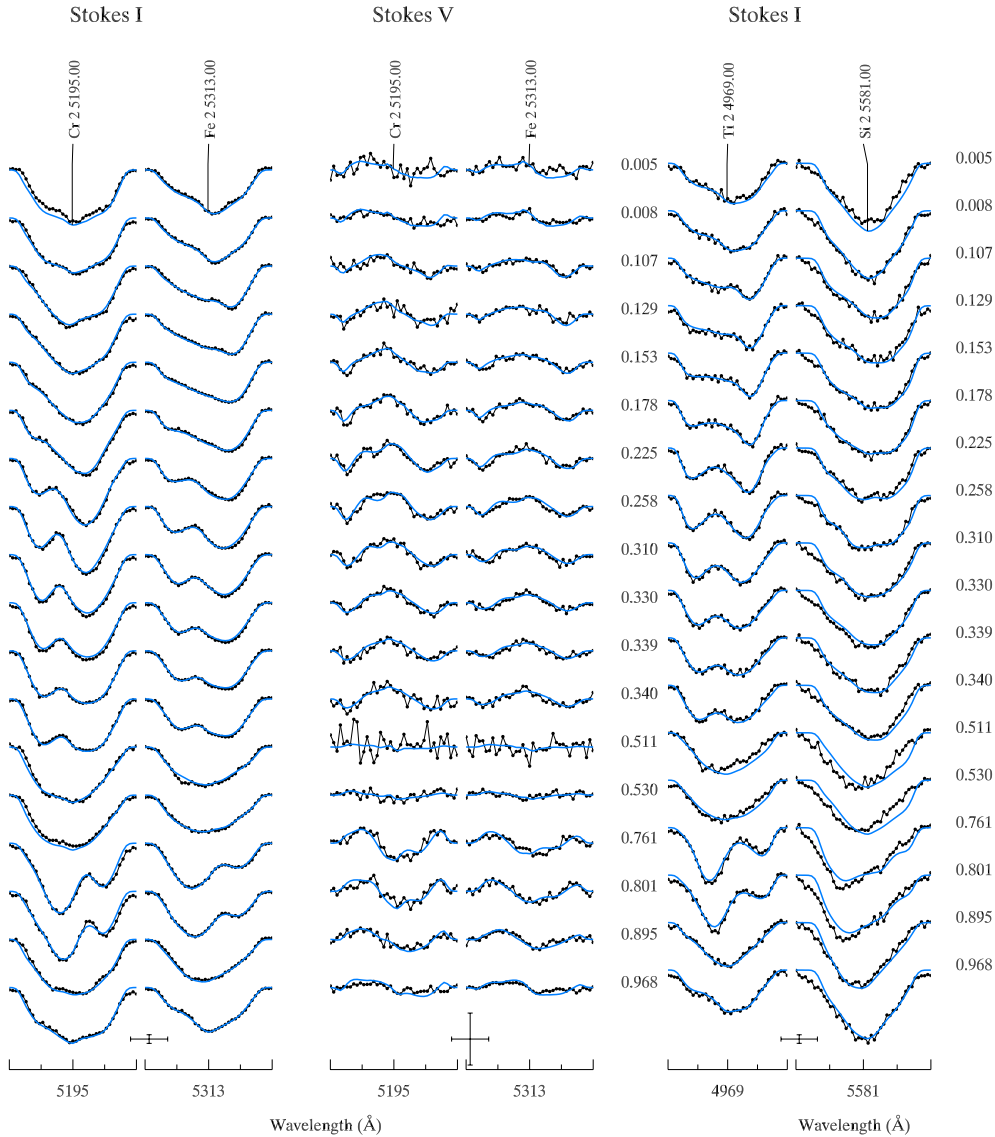


Figure 10. Fits of synthetic Stokes *I* (for Si, Ti, Cr and Fe) and Stokes *V* (for Cr and Fe) LSD profiles to observed LSD profiles. The profiles are labelled according to element, phase and mean wavelength. The bars near the bottom of the figure indicate the vertical and horizontal scale, 0.5 per cent of the continuum and 0.5 Å, respectively. Generally good fits can be seen, ranging from the high S/N Fe profiles at the best to the noisy Si profiles at the worst.

significance of these similarities and differences is unclear, as the evolution of chemical abundance spots over a star's main-sequence lifetime is not understood.

9 DISCUSSION AND CONCLUSIONS

We have analysed 20 high-resolution spectropolarimetric observations of the HD 72106 system. In these observations we see clear evidence for a magnetic field in HD 72106A. We also confirm that HD 72106B is a HAeBe star, based on emission in H α and the O I 7771 Å triplet, and that it displays no magnetic field.

There is strong evidence that the HD 72106 system is a true binary system. Both stars have the same proper motions, the same radial velocities, and *Hipparcos* solution to the system places both stars at the same distance. Although the separation of the components has remained constant in the past 90 yr, there has been a slow

systematic increase in the position angle observed. Additionally there is a wide range of possible bound orbits consistent with all available observations.

We find the age of the system to be between 6 and 13 Myr, based on the HR diagram position of the secondary (and assuming that, as a HAeBe star, the secondary is a pre-main-sequence star). Thus the system is fairly evolved, for a pre-main-sequence system. In the youngest limit, the primary would be within 1 Myr of the ZAMS. However, in the oldest limit the primary would have reached the main sequence 9 Myr ago, but it would have only passed about 1.5 per cent of its main-sequence lifetime. Thus, while it may not be on the pre-main sequence, HD 72106A is certainly very young. While it is possible that the system is not coeval and instead was formed by capture, the data available (consistent HR diagram positions and the young age of the secondary) suggest that the system truly coeval. Even with the large uncertainty in age, only a few known Ap/Bp stars approach the maximum fractional age of HD 72106A.

(Bagnulo et al. 2004; Landstreet et al. 2007). Given its evolutionary status, HD 72106A appears to represent a link between magnetic H A e Be stars and the Ap/Bp stars.

HD 72106A possesses a strong, predominantly dipolar, magnetic field. We find that a centred dipole with a polar field strength of 1230 ± 80 G and an obliquity angle of $57^\circ \pm 5^\circ$ is sufficient to provide a good fit to our observations. In HD 72106B we find no evidence of a magnetic field, with an upper limit on the longitudinal field of about 200 G. HD 72106A is one of the very youngest stars for which detailed modelling of the magnetic field geometry has been performed. The only clearly younger object is the Herbig Be star HD 200775A, which Alecian et al. (2008a) find to have an age of 0.1 ± 0.05 Myr. This star also displays a dipolar magnetic field geometry, with a magnetic field strength at the pole of 1000 ± 150 G and an obliquity of $125^\circ \pm 8^\circ$, slightly offset at $0.05 \pm 0.04 R_*$ from the star's centre (Alecian et al. 2008a). HD 200775A and HD 72106A have similar magnetic field characteristics, both of which are analogous to most main-sequence Ap/Bp stars. This provides further evidence that there is a continuum of magnetic A and B stars from the pre-main sequence through the main sequence.

We see strong chemical peculiarities in HD 72106A, particularly an underabundance of He and overabundances of Si, Ti, Cr, Fe and Nd. Thus, from the point of view of abundance analysis, HD 72106A appears to be a Bp star. This implies that chemical peculiarities can form very early in a star's lifetime, possibly even on the pre-main sequence. This conclusion is supported by the recent results of Alecian et al. (2008b) who report chemical peculiarities in the spectrum of the H A e Be star W601 in NGC 6611.

The only other star with a comparably young fractional age in which strong chemical peculiarities have been measured in detail is the main-sequence Bp star NGC 2244 334 (Bagnulo et al. 2004). This star has been on the main sequence for 2.3 ± 0.3 Myr, but is more massive and hence probably more evolved than HD 72106A, with a fractional age $\tau = 0.02 \pm 0.01$ (Hensberge, Pavlovski & Verschuere 2000; Bagnulo et al. 2004). Bagnulo et al. (2004) find that the star has a temperature of $15\,000 \pm 1000$ K and an observed longitudinal magnetic field strength of 9000 G. The authors find a strong underabundance of He by ~ 1 dex, as well as strong overabundances of Si and Fe by ~ 1 dex, and Ti and Cr by ~ 2 dex. These peculiarities are similar to those of HD 72106A, suggesting that there are important similarities between these stars.

We find clear evidence for surface abundance inhomogeneities in HD 72106A. MDI suggests that there is a large spot of overabundance in Si, Ti, Cr and Fe near the positive magnetic pole. These inhomogeneities are similar to those seen in Ap/Bp stars, further supporting a link between magnetic H A e Be stars and Ap/Bp stars. The Doppler reconstructions presented here are the earliest stage of intermediate-mass stellar evolution ever mapped.

HD 72106B, in contrast to HD 72106A, appears to be chemically normal. While this is not the first binary system containing a magnetic chemically peculiar and a chemically normal star (e.g. Carrier et al. 2002), it does raise the question of how such systems are produced. Being a young binary, it is likely that the stars of HD 72106 formed at the same time from approximately the same material. The temperatures of these stars differ by only 2250 ± 1120 K, and the $v \sin i$ values of the stars differ by only ~ 10 km s $^{-1}$. However, one star is magnetic and chemically peculiar while the other is not. This suggests that, whatever mechanism gave rise to the difference between the stars, it must be rather sensitive to the particulars of the star's initial conditions.

It is instructive to compare the characteristics of HD 72106A to those of the magnetic H A e Be stars HD 104237 and HD 190073.

HD 104237 was observed to possess a magnetic field by Donati et al. (1997) and confirmed by Donati (2000) and Alecian et al. (2008c), with a longitudinal field strength of ~ 50 G (Donati, private communication). HD 104237 has a mass of about $2.3 M_\odot$ and an age of about 2 Myr (van den Ancker, de Winter & Tjin A Djie 1998). Acke & Waelkens (2004) performed an abundance analysis of HD 104237 using equivalent widths, and found approximately solar abundances for a range of elements, including Si, Cr and Fe. The star HD 190073 was reported to possess a magnetic field by Catala et al. (2007), with a longitudinal magnetic field strength of 74 ± 10 G. Catala et al. (2007) derive a mass of $2.85 \pm 0.25 M_\odot$ and an age of 1.2 ± 0.6 Myr (measured from the birth line) for this star. Acke & Waelkens (2004) also studied the surface chemistry of this star and found roughly solar abundances.

Thus it appears that the majority of known magnetic H A e Be stars are chemically normal, though some peculiar stars seem to exist (such as NGC 6611 W601, Alecian et al. 2008b). This is in contrast to main-sequence Ap/Bp stars, in which magnetic fields are nearly always found with chemical peculiarities. More analysis of chemical abundances in magnetic H A e Be stars must be performed, with an eye to identifying chemically peculiar objects. Interestingly, the young chemically peculiar stars HD 72106A and NGC 2244 334 display no emission, while HD 104237 and HD 190073 both display significant emission in their spectra. This suggests that HD 104237 and HD 190073 may still be undergoing significant accretion or mass loss while HD 72106A and NGC 2244 334 are not. Thus it may be that accretion or mass loss mixes the stellar atmosphere, inhibiting the build-up of chemical peculiarities through diffusion. However, once accretion halts chemical peculiarities may arise quickly.

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