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Gravitational Collapse and Star Formation

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The Cosmic Matter Cycle





Thackeray's Globules in IC 2944





NASA and The Hubble Heritage Team (STScI/AURA) • Hubble Space Telescope WFPC2 • STScI-PRC02-01



Is a cloud stable or will it tend to collapse due to gravitational forces ?

1D, gravity included:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) = 0$$

$$\frac{\partial}{\partial t}(\rho u) + \frac{\partial}{\partial x}(\rho u u) = -\frac{\partial P}{\partial x} - \rho \frac{\partial \Phi}{\partial x}$$

$$\frac{\partial^2 \Phi}{\partial x^2} = 4\pi G\rho$$

$$P = c_s^2 \rho$$

isothermal gas

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$$P = c_s^2 \rho$$

$$P = P_0 + P_1(x, t)$$

$$\rho = \rho_0 + \rho_1(x, t)$$

$$u = u_1(x, t)$$

$$\Phi = \Phi_0 + \Phi_1(x, t)$$

 $\frac{\partial \rho_1}{\partial t}$

linearised equations:

isothermal perturbation

$$+ \rho_0 \frac{\partial u_1}{\partial x} = 0 \frac{\partial u_1}{\partial t} = - \frac{c_s^2}{\rho_0} \frac{\partial \rho_1}{\partial x} - \frac{\partial \Phi_1}{\partial x} \frac{\partial^2 \Phi_1}{\partial x^2} = 4\pi G \rho_1 P_1 = c_s^2 \rho_1 P = P_0 + P_1(x,t) \rho = \rho_0 + \rho_1(x,t) u = u_1(x,t) \Phi = \Phi_0 + \Phi_1(x,t)$$

 $\frac{\partial \rho_1}{\partial t}$

linearised equations:

linear homogeneous system of differential equations with constant coefficients for ρ_1 , u_1 and Φ_1

isothermal perturbation

small-amplitude disturbances in gas which initially is at rest with constant pressure and density

 $\rho_{I}, u_{I}, \Phi_{I} \propto exp[i(kx + \omega t)]$

$$+ \rho_0 \frac{\partial u_1}{\partial x} = 0 \frac{\partial u_1}{\partial t} = - \frac{c_s^2}{\rho_0} \frac{\partial \rho_1}{\partial x} - \frac{\partial \Phi_1}{\partial x} \frac{\partial^2 \Phi_1}{\partial x^2} = 4\pi G \rho_1 P_1 = c_s^2 \rho_1 P = P_0 + P_1(x, t) \rho = \rho_0 + \rho_1(x, t) u = u_1(x, t) \Phi = \Phi_0 + \Phi_1(x, t)$$

 $i \omega \rho_1$

linearised equations:

isothermal perturbation

$$\rho_{i}, u_{i}, \Phi_{i} \propto exp[i(kx + \omega t)]$$

$$+ \rho_0 i k u_1 = 0
i \omega u_1 = -\frac{c_s^2}{\rho_0} i k \rho_1 - i k \Phi_1
-k^2 \Phi_1 = 4 \pi G \rho_1
P_1 = c_s^2 \rho_1
P = P_0 + P_1(x, t)
\rho = \rho_0 + \rho_1(x, t)
u = u_1(x, t)
\Phi = \Phi_0 + \Phi_1(x, t)$$

linearised equations:

isothermal perturbation

$$\rho_{i}, u_{i}, \Phi_{i} \propto exp[i(kx + \omega t)]$$

$$\begin{split} \omega \rho_{1} + \rho_{0} k u_{1} &= 0 \\ \omega u_{1} &= -\frac{c_{s}^{2}}{\rho_{0}} k \rho_{1} - k \Phi_{1} \\ -k^{2} \Phi_{1} &= 4 \pi G \rho_{1} \\ P_{1} &= c_{s}^{2} \rho_{1} \\ P_{1} &= c_{s}^{2} \rho_{1} \\ P_{1} &= c_{s}^{2} \rho_{1} \\ p &= \rho_{0} + \rho_{1}(x, t) \\ \mu &= u_{1}(x, t) \\ J &= \Phi_{0} + \Phi_{1}(x, t) \end{split}$$

linearised equations:

Actions: $\omega \rho_1 + \rho_0 k u_1 = 0$ $\frac{c_s^2}{\rho_0} k \rho_1 + \omega \quad u_1 + k \quad \Phi_1 = 0$ $4\pi G \rho_1 + k^2 \Phi_1 = 0$

isothermal perturbation

$$\rho_{i}, u_{i}, \Phi_{i} \propto exp[i(kx + \omega t)]$$

+
$$k^{2} \Phi_{1} = 0$$

 $P_{1} = c_{s}^{2} \rho_{1}$
 $P = P_{0} + P_{1}(x, t)$
 $\rho = \rho_{0} + \rho_{1}(x, t)$
 $u = u_{1}(x, t)$
 $\Phi = \Phi_{0} + \Phi_{1}(x, t)$

dispersion relation:

$$\omega^2 = k^2 c_s^2 - 4\pi G \rho_0$$

 $k^2 c_s^2 - 4\pi G \rho_0 > 0 \implies \omega \text{ real}$ (periodic, stable) $k^2 c_s^2 - 4\pi G \rho_0 < 0 \implies \omega \text{ imaginary (growing perturb.)}$

isothermal perturbation

$$\rho_{i'} u_{i'} \Phi_{i} \propto \exp[i(kx + \omega t)]$$

$$P_1 = c_s^2 \rho_1$$

$$P = P_0 + P_1(x, t)$$

$$\rho = \rho_0 + \rho_1(x, t)$$

$$u = u_1(x, t)$$

$$\Phi = \Phi_0 + \Phi_1(x, t)$$

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critical wavenumber:

$$k_{J} = \left(\frac{4\pi \ G \ \rho_{0}}{c_{s}^{2}}\right)^{1/2}$$
$$\lambda_{J} = c_{s} \left(\frac{\pi}{G \ \rho_{0}}\right)^{1/2} \qquad (\lambda = 2\pi/k)$$

or critical wavelength:

Jeans criterion: $\lambda_J = c_s \left(\frac{\pi}{G \rho_0}\right)^{1/2}$ $c_s = \left(\frac{P_0}{\rho_0}\right)^{1/2}$

perturbations with $\lambda > \lambda_J$ are unstable (growing amplitudes) leading to collapse due to gravity

Jeans mass:

$$M_{J} = \rho_{0} \left(\frac{\lambda_{J}}{2}\right)^{3}$$

$$M_{J} = \rho_{0} \left(\frac{\pi \ k \ T_{0}}{4 \ \mu \ m_{u} \ G \ \rho_{0}}\right)^{3/2} \propto T_{0}^{3/2} \ \rho_{0}^{-1/2}$$

or



Is a cloud stable or will it tend to collapse due to gravitational forces ?

Jeans criterion:

given density, temperature

 $\downarrow critical mass M_{J}$ associated with a critical length scale for a disturbance

Table 19.1 Typical States of Gas in the Interstellar Medium

State of Gas	Primary Constituent	Approximate Temperature	Approximate Density (atoms per cm ³)	Description
Hot bubbles	Ionized hydrogen	1,000,000 K	0.01	Pockets of gas heated by supernova shock waves
Warm atomic gas	Atomic hydrogen	10,000 K	1	Fills much of galactic disk
Cool atomic clouds	Atomic hydrogen	100 K	100	Intermediate stage of star–gas–star cycle
Molecular clouds	Molecular hydrogen	30 K	300	Regions of star formation
Molecular cloud cores	Molecular hydrogen	60 K	10,000	Star-forming clouds

Star Formation in Dense Clouds





Star Formation in Dense Clouds



Gaseous Pillars · M16

PRC95-44a · ST Scl OPO · November 2, 1995 J. Hester and P. Scowen (AZ State Univ.), NASA

Star Formation in Dense Clouds









Is a cloud stable or will it tend to collapse due to gravitational forces ?

Jeans criterion:

given density, temperature

↓ critical mass M_J associated with a critical length scale for a disturbance

How fast will the cloud collapse ?

Typical Time Scale of the Collapse

Spherical cloud collapsing under its own gravity:

motion of mass shells determined by gravity (free fall), assuming that gas pressure and other forces can be neglected

 \Rightarrow free fall time

$$t_{ff} = \left(\frac{3 \pi}{32 G \rho_0}\right)^{1/2}$$

 $\rho_0 \dots$ initial (mean) density of the cloud $\rho_0 = M_{cloud} / (4\pi R_{cloud}^3 / 3)$ where R_{cloud} is the initial radius of the cloud and M_{cloud} its total mass

Typical Time Scale of the Collapse







... but reality is more complicated than that:

- influence of nearby stars (radiation, winds)
- rotation & turbulence on various scales
- young stellar objects: disks and jets ...



simulations by Banerjee et al.(2004): density in the disk plane and perpendicular to it



simulations by Banerjee et al.(2004): density in the disk plane and perpendicular to it

Young Stellar Objects: Disks & Jets





A jet model for the formation of Herbig–Haro objects. The dashed curve is the contact discontinuity which separates shocked cloud gas from shocked jet gas.

Young Stellar Objects: Disks & Jets



Jets from Young Stars

PRC95-24a · ST Scl OPO · June 6, 1995 C. Burrows (ST Scl), J. Hester (AZ State U.), J. Morse (ST Scl), NASA